



R&DE (Engineers), DRDO

2D Theory of Elasticity

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Summary

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- Field equations
- Boundary conditions
- Problem formulation
- Plane strain
- Plane stress
- Airy's stress function
- Axially loaded bar
- Pure bending of beam



Introduction

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- In principle all practical problems – 3D problems – very complex – difficult to handle
- Reasonable assumptions – bring down the complexity in the problem – reduces a dimension => 2D problem
- 2D problems – two independent variables – static problem – three independent variables – dynamic problems
- Field variables, $\phi = \phi(x, y, t)$



Introduction

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- Field variable is independent of third dimension $\Rightarrow \frac{\partial \phi}{\partial z} = 0$
- Continuum approach – the following laws hold good
 - Conservation of mass
 - Principle of momentum
 - Principle of moment of momentum
 - First law of thermodynamics
 - Second law of thermodynamics - inequality



Introduction

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- No inertia effects – quasi-static / static
- Small deformation and rotations
- Material \Rightarrow isotropic – same properties in all the direction
- Homogeneous – material properties are uniform over the domain
- Linear elastic – Hooke's law holds good \Rightarrow Hookean material



Field equations

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■ Strain-displacement equations

$$\varepsilon_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) = \frac{1}{2} (u_{i,j} + u_{j,i})$$

Six equation

■ Compatibility equations

$$\varepsilon_{ij,kl} + \varepsilon_{kl,ij} - \varepsilon_{ik,jl} - \varepsilon_{jl,ik} = 0$$

Six equations – three independent



Field equations

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- Equilibrium equations –

$$\sigma_{ij,j} + B_i = 0$$

Three equations

- Constitutive law – Linear elastic – Hooke's law

$$\sigma_{ij} = \lambda \varepsilon_{kk} \delta_{ij} + 2 \mu \varepsilon_{ij}$$

$$\varepsilon_{ij} = \frac{1 + \nu}{E} \sigma_{ij} - \frac{\nu}{E} \sigma_{kk} \delta_{ij}$$

Six equations

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Field equations

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- Total – 15 equations = 3 Equilibrium + 6 strain displacement + 6 constitutive laws
- Unknowns => 3 Displacements, 6 stress components and 6 strain components
- Boundary value problem casted as

$$\mathfrak{I} \left(\underbrace{u_i, \varepsilon_{ij}, \sigma_{ij}}_{\text{Unknowns}}, \underbrace{\lambda, \mu}_{\text{Lame's constants}}, B_i \right) = 0$$

Body load

Not easy to solve analytically 15 equations to find out 15 unknowns

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Boundary conditions

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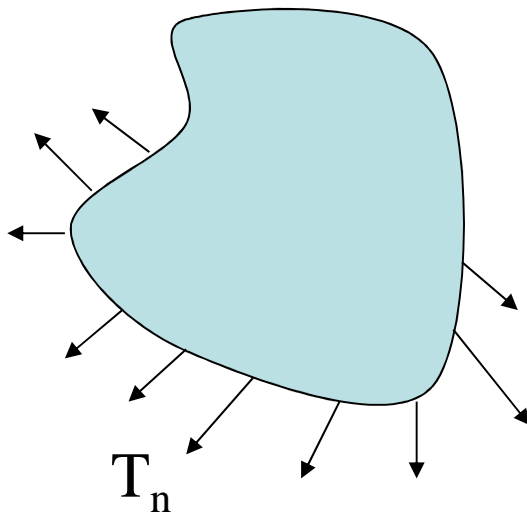
- Solution of a differential (partial) equations – boundary conditions required
- Boundary conditions influence the solution – field variable in the domain
- BCs play a very essential role for properly formulating and solving elasticity problems
- In elasticity – boundary conditions – displacements, tractions and combination



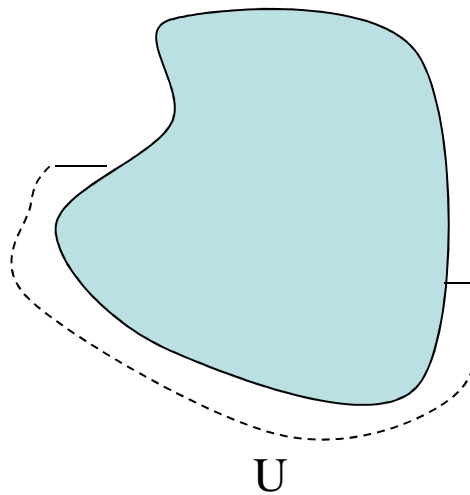
Boundary conditions

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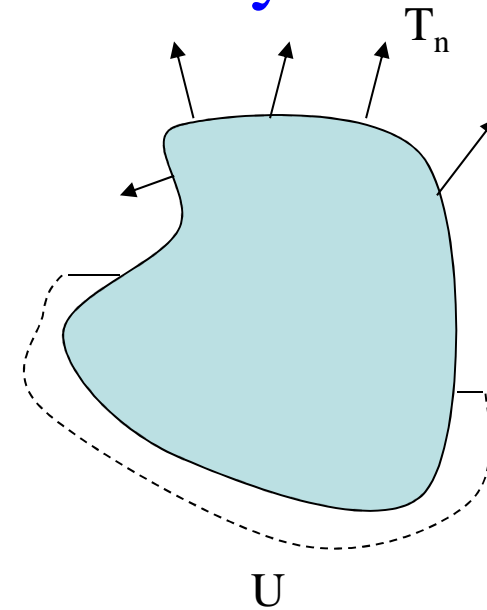
Various boundary conditions in elasticity –



Traction boundary conditions



Displacement boundary conditions



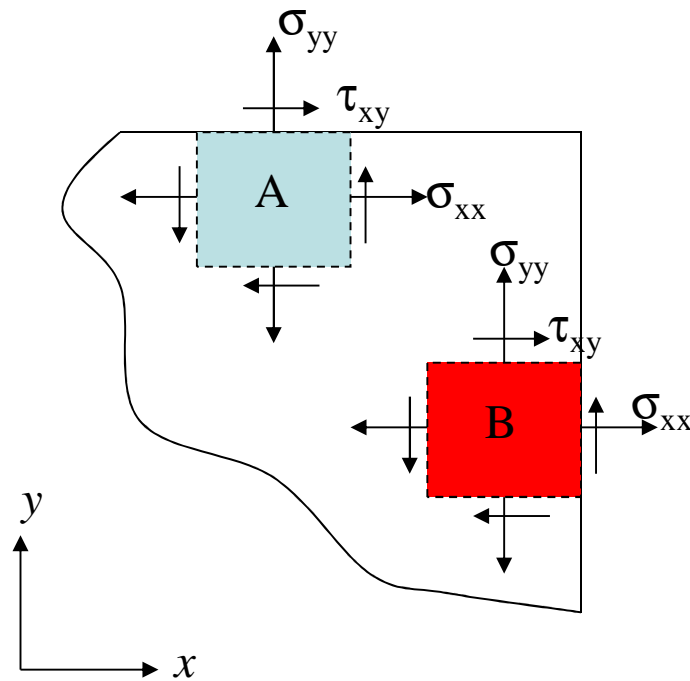
Mixed boundary conditions



Boundary conditions

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■ Traction –



$$T_x = l \sigma_{xx} + m \tau_{xy}$$

$$T_y = l \tau_{xy} + m \sigma_{yy}$$

At 'A' – Direction cosines of area normal $\Rightarrow (0, 1, 0)$

σ_{xx} is inside the boundary – this need not be zero

$$T_x = \tau_{xy}; \quad T_y = \sigma_{yy}$$

At 'B' – Direction cosines of area normal $\Rightarrow (1, 0, 0)$

$$T_x = \sigma_{xx}; \quad T_y = \tau_{xy}$$

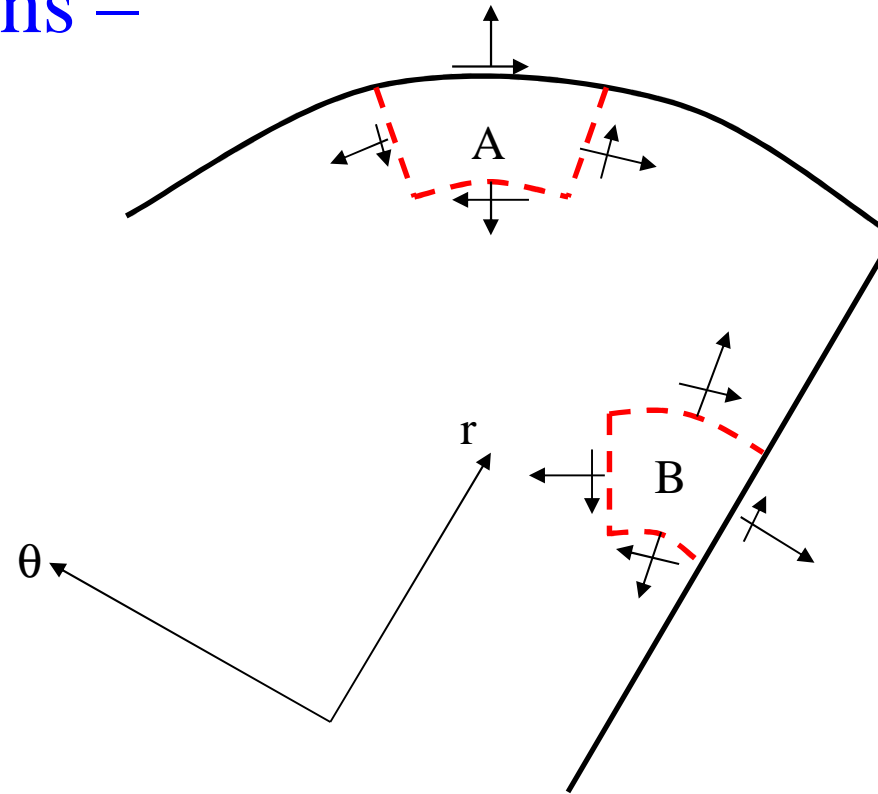
σ_{yy} is inside the boundary – this need not be zero



Boundary conditions

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■ Traction –



Traction in Polar co-ordinate system

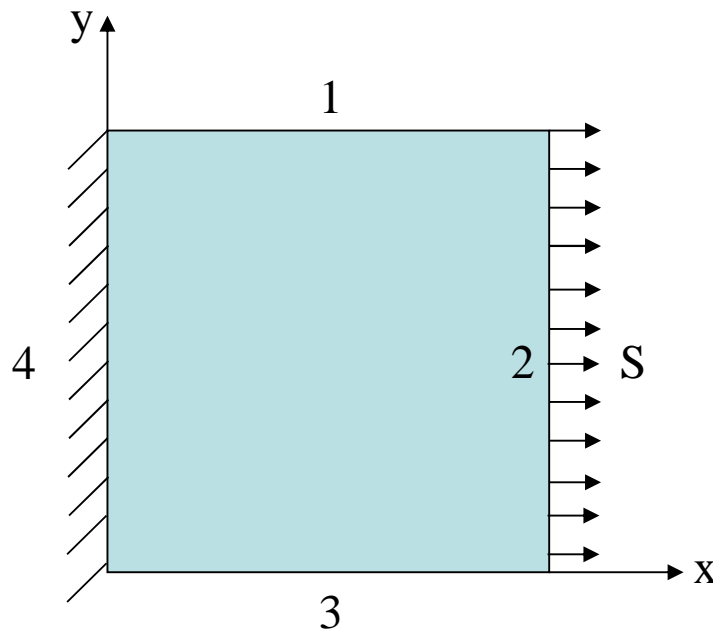
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Boundary conditions

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■ Traction –



Boundary – 4

$$u = 0; v = 0$$

Boundary – 1 – traction free

$$DC \Rightarrow (0, 1, 0)$$

$$T_x = 0. \sigma_{xx} + 1. \tau_{xy} = 0$$

$$\Rightarrow T_x = \tau_{xy} = 0$$

$$\Rightarrow T_y = 0. \tau_{xy} + 1. \sigma_{yy} = 0 \Rightarrow \sigma_{yy} = 0$$

Boundary – 2 – DCs (1, 0, 0)

$$T_x = 1. \sigma_{xx} + 0. \tau_{xy} = S = \sigma_{xx}$$

$$T_y = 1. \tau_{xy} + 0. \sigma_{yy} = 0 = \tau_{xy}$$

Boundary – 3 – DCs (0, -1, 0)

$$T_x = 0. \sigma_{xx} - 1. \tau_{xy} = 0$$

$$T_x = \tau_{xy} = 0$$

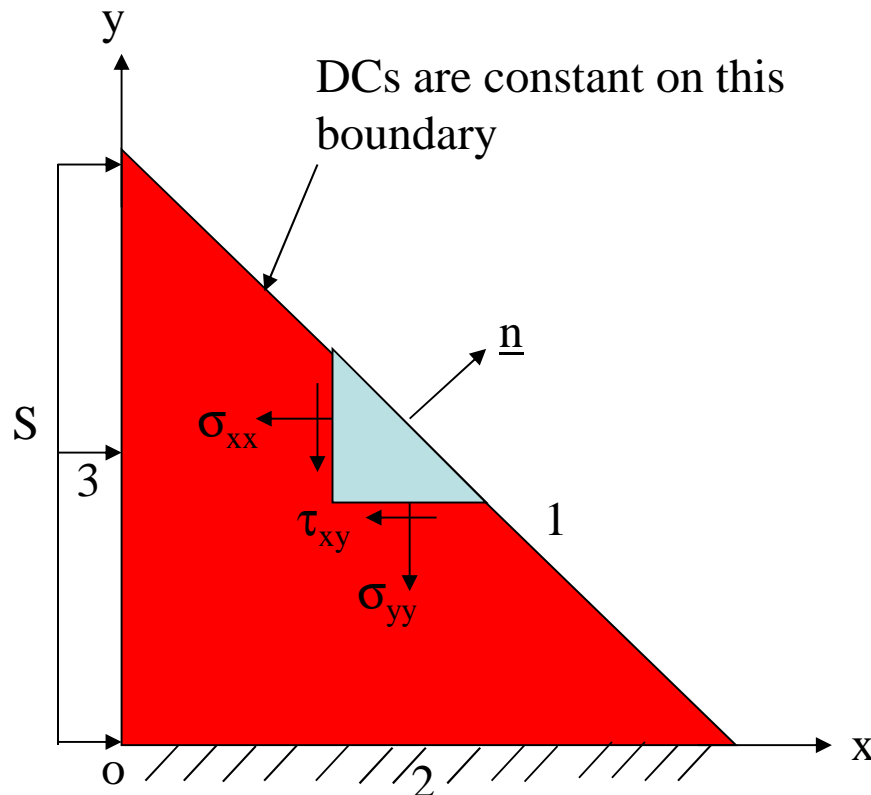
$$T_y = 0. \tau_{xy} - 1. \sigma_{yy} = 0 \Rightarrow \sigma_{yy} = 0$$



Boundary conditions

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■ Traction



Individual stress components are not zero on boundary - 1

Boundary - 1 – DCs ($l, m, 0$)

This is a traction free boundary

$$T_x = l\sigma_{xx} + m\tau_{xy} = 0$$

$$T_y = m\sigma_{xx} + l\tau_{xy} = 0$$

Boundary - 2 – Constrained boundary – $U, V = 0$

Boundary - 3 – Traction is acting

DCs ($-1, 0, 0$)

$$T_x = -l\sigma_{xx} + m\tau_{xy} = S \Rightarrow \sigma_{xx} = -S$$

$$T_y = m\sigma_{xx} + l\tau_{xy} = 0 \Rightarrow \tau_{xy} = 0$$



Problem formulation

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- 3D elasticity 15 equations and 15 unknowns - 2D elasticity – 8 equations and 8 unknowns
- Reformulating elasticity problems – mathematically convenient way
- Two approaches –
 - Displacement based
 - Stress based formulations
- Displacement based – express field equation in terms of displacements, u_i – Navier equation
- Stress based – express field equation in terms of stresses, σ_{ij}



Problem formulation

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- In displacement based formulation
 - Primary variable – displacement – solution is obtained for displacements – stresses and strains secondary / derived variables
 - Strains – derivation of displacements
 - Stresses – use constitutive law
- In stress based formulation
 - Primary variable – stresses – solution is obtained for stresses – strains and displacements secondary variables
 - Strain compatibility equations



Displacement formulation

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- Governing equation expressed in displacements
- Field equations –

$$\text{Equilibrium} \quad \sigma_{ij,j} + B_i = 0 \quad - (1)$$

$$\text{Strain - displacement} \quad \varepsilon_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) \quad - (2)$$

$$\text{Constitutive law} \quad \sigma_{ij} = \lambda \varepsilon_{kk} \delta_{ij} + 2\mu \varepsilon_{ij} \quad - (3)$$



Displacement formulation

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- Differentiate (3) wrt 'j'

$$\sigma_{ij,j} = \lambda \varepsilon_{kk,j} \delta_{ij} + 2\mu \varepsilon_{ij,j} \quad - (4)$$

Plug this in (1)

$$\sigma_{ij,j} + B_i = \lambda \varepsilon_{kk,j} \delta_{ij} + 2\mu \varepsilon_{ij,j} + B_i = 0 \quad - (5)$$

Express all strain components in (5) in terms of displacements – Use equation (2)

$$\lambda u_{k,kj} \delta_{ij} + \mu (u_{i,jj} + u_{j,ij}) + B_i = 0$$

$$\lambda u_{k,ki} + \mu u_{j,ji} + \mu u_{i,jj} + B_i = 0$$



Displacement formulation

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$$\lambda u_{k,ki} + \mu u_{j,ji} + \mu u_{i,jj} + B_i = 0$$

$$(\lambda + \mu) u_{j,ji} + \mu u_{i,jj} + B_i = 0$$

This gives three equations for $i = 1, 2$ and 3

For $i = 1 \Rightarrow x$ - direction

$$(\lambda + \mu) \frac{\partial}{\partial x_1} \left(\frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3} \right) + \mu \left(\frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} + \frac{\partial^2}{\partial x_3^2} \right) u_1 + B_1 = 0$$

Combining all three equations into a single equation

$$(\lambda + \mu) \nabla (\nabla \cdot \underline{\underline{U}}) + \mu \nabla^2 \underline{\underline{U}} + \underline{\underline{B}} = 0$$



Displacement formulation

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- Governing equations in terms of displacements
– three equations - reduced from '15' equations

$$(\lambda + \mu)u_{j,jj} + \mu u_{i,jj} + B_i = 0 \quad - (U)$$

$$\mathfrak{I}(u_i, \lambda, \mu, B_i) = 0 \quad \text{Navier equation}$$

Equations (U) known as Navier's or Lamé's equations

Boundary conditions $\Rightarrow u_i = U_i$ on 'S' boundary $u_i = u_i(x_1, x_2, x_3)$

Solve second order PDEs – get displacements

Use strain – displacement relations for calculating strains

Constitutive relations for stresses



Stress formulation

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- In displacement formulation compatibility equations play no role – displacement obtained are single valued and continuous
- In stress formulation express governing equation in terms of stresses
- Stresses – primary variables – obtain secondary variables displacements from strain-displacement relations
- Integrate ε - U relations – compatibility equations come into play



Stress formulation

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- Boundary conditions – only tractions
- From the field equations eliminate strains and displacements
- Compatibility equation

$$\epsilon_{ij,kl} + \epsilon_{kl,ij} - \epsilon_{ik,jl} - \epsilon_{jl,ik} = 0 \quad - \quad (1)$$

For $k = l$ first three compatibility equations are obtained

Constitutive law $\epsilon_{ij} = \frac{1+\nu}{E} \sigma_{ij} - \frac{\nu}{E} \sigma_{kk} \delta_{ij}$

Differentiate this wrt k and l , plug in (1)

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Stress formulation

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- Three compatibility equations for $k=l = 1, 2$ and 3

$$\frac{\partial^2 \epsilon_{xx}}{\partial y^2} + \frac{\partial^2 \epsilon_{yy}}{\partial x^2} = \frac{\partial^2 \gamma_{xy}}{\partial x \partial y}$$
$$(i = j = 1, k = l = 2)$$

$$\frac{\partial^2 \epsilon_{yy}}{\partial z^2} + \frac{\partial^2 \epsilon_{zz}}{\partial y^2} = \frac{\partial^2 \gamma_{yz}}{\partial y \partial z}$$
$$(i = j = 2, k = l = 3)$$

$$\frac{\partial^2 \epsilon_{xx}}{\partial z^2} + \frac{\partial^2 \epsilon_{zz}}{\partial x^2} = \frac{\partial^2 \gamma_{xz}}{\partial x \partial z}$$
$$(i = j = 3, k = l = 1)$$



Stress formulation

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This gives –

$$\sigma_{ij,kk} + \sigma_{kk,ij} - \sigma_{ik,jk} - \sigma_{jk,ik} =$$
$$\frac{\nu}{1+\nu} \left(\sigma_{mm,kk} \delta_{ij} + \sigma_{mm,ij} \delta_{kk} - \sigma_{mm,jk} \delta_{ik} - \sigma_{mm,ik} \delta_{jk} \right)$$

$$\delta_{kk} = 3$$

$$\text{equilibrium eqn. } \sigma_{ij,j} + B_i = 0$$

$$\sigma_{ij,j} = -B_i$$

Make use of equilibrium equation to reduce further



Stress formulation

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Governing equation in terms of stress –

$$\sigma_{ij, kk} + \frac{1}{1 + \nu} \sigma_{kk, ij} = \frac{\nu}{1 + \nu} \sigma_{mm, kk} \delta_{ij} - B_{i, j} - B_{j, i}$$

$$\mathfrak{I}(\sigma_{ij}, \lambda, \mu, B_i) = 0$$

Compatibility equation in terms of stresses, known as

“Beltrami-Michell” compatibility equation

Six equations can be obtained for $i, j = 1, 2, 3$

Only three are independent – similar to strain compatibility equations

Three more equations are required – complemented by three equilibrium equations

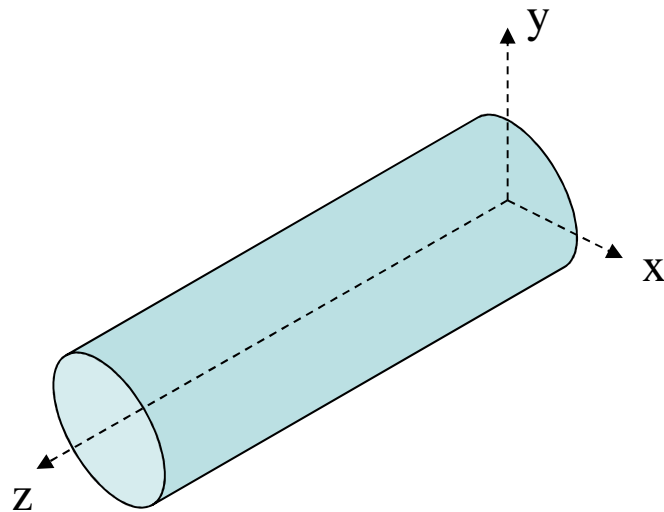
- Six unknowns – six equations
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Plane strain

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- Out-of-plane dimensions are large compared to in-plane dimensions



Reduce a 3D problem to 2D problem

All loads and boundary conditions independent of 'z' co-ordinate

The deformation in 'z' direction assumed to be zero. $w = 0$

Deformations in 'x' and 'y' directions depend on (x, y) only

$$u = u(x, y) \text{ and } v = v(x, y)$$

All cross-sections have same displacements

Out of plane strains vanish



Plane strain

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- In plane strain $\varepsilon_{zz} = 0$, $\gamma_{yz} = 0$ and $\gamma_{zx} = 0$
- Linear elastic material – Hooke's law

$$\sigma_{ij} = \lambda \varepsilon_{kk} \delta_{ij} + 2\mu \varepsilon_{ij}$$

$$\sigma_{xx} = \lambda (\varepsilon_{xx} + \varepsilon_{yy}) + 2\mu \varepsilon_{xx}$$

$$\sigma_{yy} = \lambda (\varepsilon_{xx} + \varepsilon_{yy}) + 2\mu \varepsilon_{yy}$$

$$\sigma_{zz} = \lambda (\varepsilon_{xx} + \varepsilon_{yy})$$

$$\tau_{xy} = 2\mu \varepsilon_{xy} = \mu \gamma_{xy}$$



Plane strain

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■ Equilibrium equations –

$$\sigma_{ij,j} + B_i = 0$$

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + B_x = 0 \quad - \quad (1)$$

$$\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{xx}}{\partial y} + B_y = 0 \quad - \quad (2)$$

Use constitutive law and strain-displacement equations in (1)

$$\frac{\partial}{\partial x} \left(\lambda (\epsilon_{xx} + \epsilon_{yy}) + 2\mu \epsilon_{xx} \right) + \frac{\partial}{\partial y} (\mu \gamma_{xy}) + B_i = 0$$

$$\frac{\partial}{\partial x} \left[\lambda \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + 2\mu \frac{\partial u}{\partial x} \right] + \frac{\partial}{\partial y} \left[\mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right] + B_x = 0$$



Plane strain

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Governing equation in x – direction

$$\mu \nabla^2 u + (\lambda + \mu) \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + B_x = 0$$

Use second equilibrium equation – convert to displacements

$$\mu \nabla^2 v + (\lambda + \mu) \frac{\partial}{\partial y} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + B_y = 0$$

In each of the above two equations, equilibrium, constitutive law and displacement – strain equations are inbuilt - ***Displacement formulation***

Eight equations => reduced to two equations in displacements

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Plane strain

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■ Stress based formulation -

In 2D case, compatibility equation
$$\frac{\partial^2 \epsilon_{xx}}{\partial y^2} + \frac{\partial^2 \epsilon_{yy}}{\partial x^2} = \frac{\partial^2 \gamma_{xy}}{\partial x \partial y}$$

Use constitutive law – convert into stresses

$$\epsilon_{xx} = \frac{\sigma_{xx}}{E} - \frac{\nu}{E} (\sigma_{yy} + \sigma_{zz})$$

$$\epsilon_{yy} = \frac{\sigma_{yy}}{E} - \frac{\nu}{E} (\sigma_{xx} + \sigma_{zz})$$

$$\epsilon_{zz} = 0 = \frac{\sigma_{zz}}{E} - \frac{\nu}{E} (\sigma_{xx} + \sigma_{yy}) \Rightarrow \sigma_{zz} = \nu (\sigma_{xx} + \sigma_{yy})$$

$$\gamma_{xy} = \frac{2(1+\nu)}{E} \tau_{xy}$$

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Plane strain

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■ Compatibility equation in terms of stresses

$$\frac{\partial^2}{\partial x^2} [(1-\nu)\sigma_{yy} - \nu\sigma_{xx}] + \frac{\partial^2}{\partial y^2} [(1-\nu)\sigma_{xx} - \nu\sigma_{yy}] = 2 \frac{\partial^2 \tau_{xy}}{\partial x \partial y}$$

Above equation is obtained from compatibility, constitutive equations

Use equilibrium equations to eliminate shear stress from above equation

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + B_x = 0 \Rightarrow \frac{\partial}{\partial x} \left(\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + B_x \right) = 0$$

$$\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{xx}}{\partial y} + B_y = 0 \Rightarrow \frac{\partial}{\partial y} \left(\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{xx}}{\partial y} + B_y \right) = 0$$

Add

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Plane strain

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This is obtained as –

$$\frac{\partial^2 \sigma_{xx}}{\partial x^2} + \frac{\partial^2 \sigma_{yy}}{\partial y^2} + 2 \frac{\partial^2 \tau_{xy}}{\partial x \partial y} + \frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} = 0$$

$$2 \frac{\partial^2 \tau_{xy}}{\partial x \partial y} = - \left(\frac{\partial^2 \sigma_{xx}}{\partial x^2} + \frac{\partial^2 \sigma_{yy}}{\partial y^2} + \frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} \right)$$

$$\frac{\partial^2}{\partial x^2} [(1-\nu)\sigma_{yy} - \nu\sigma_{xx}] + \frac{\partial^2}{\partial y^2} [(1-\nu)\sigma_{xx} - \nu\sigma_{yy}] = - \left(\frac{\partial^2 \sigma_{xx}}{\partial x^2} + \frac{\partial^2 \sigma_{yy}}{\partial y^2} + \frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} \right)$$

Simplifying this expression

$$\nabla^2 (\sigma_{xx} + \sigma_{yy}) = - \frac{1}{1-\nu} \left(\frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} \right)$$

Single governing equation for stress based formulation of plane strain problems



Plane strain

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■ Governing equation

$$\nabla^2 (\sigma_{xx} + \sigma_{yy}) = -\frac{1}{1-\nu} \left(\frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} \right)$$

In this equation, equilibrium, constitutive and compatibility equations are inbuilt.

Three stress components in 2D – σ_{xx} , σ_{yy} , τ_{xy}

Only one governing equation – three unknowns

In 2D two stress equilibrium equations exist – use these two with above governing equation

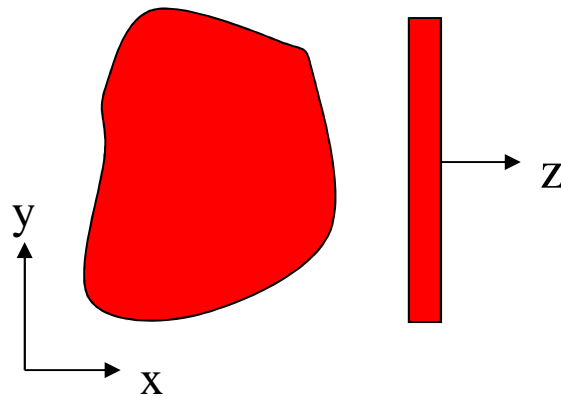
Boundary conditions given in terms of tractions. Relate them to stresses



Plane stress

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- Third dimension – much smaller than in-plane dimensions



Thin plates – structures

Reducing a 3D problem to a 2D problem

Out of plane stresses σ_{zz} , τ_{xz} and τ_{yz} negligible

Body can have out of plane deformation

ϵ_{zz} is non-zero, γ_{xz} and $\gamma_{yz} = 0$

Stresses are functions of in-plane displacements

$$\sigma_{xx} = \sigma_{xx}(x, y),$$

$$\sigma_{yy} = \sigma_{yy}(x, y),$$

$$\tau_{xy} = \tau_{xy}(x, y)$$

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Plane stress

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- Formulation of plane stress problem – two approaches
 - Displacement based
 - Stress based approach
- Formulation procedure – same as in plane strain
- Formulation takes care of non-zero out of normal strain and zero normal stress
- Use field equations



Plane stress

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■ Displacement based approach –

$$\begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix} = \frac{E}{(1-\nu^2)} \begin{bmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & \frac{(1-\nu)}{2} \end{bmatrix} \begin{Bmatrix} \epsilon_x \\ \epsilon_y \\ \gamma_{xy} \end{Bmatrix} \quad \text{Plug in stress equilibrium equations}$$

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + B_x = 0; \quad \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{xx}}{\partial y} + B_y = 0$$

$$\frac{\partial}{\partial x} \left[\frac{E}{1-\nu^2} \epsilon_{xx} + \frac{E\nu}{1-\nu^2} \epsilon_{yy} \right] + \frac{\partial}{\partial y} \left[\frac{E}{2(1+\nu)} \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right] + B_x = 0$$

$$\frac{\partial}{\partial y} \left[\frac{E}{1-\nu^2} \epsilon_{yy} + \frac{E\nu}{1-\nu^2} \epsilon_{xx} \right] + \frac{\partial}{\partial x} \left[\frac{E}{2(1+\nu)} \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right] + B_y = 0$$

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Plane stress

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- Use strain displacement equations to convert strains into displacements
- Governing equations in terms of displacements

$$\mu \nabla^2 u + \frac{E}{2(1-\nu)} \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + B_x = 0$$

$$\mu \nabla^2 v + \frac{E}{2(1-\nu)} \frac{\partial}{\partial y} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + B_y = 0$$

In the above two equations – equilibrium, constitutive laws and strain-displacements equations



Plane stress

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- Stress based approach – field equations in plane stress problems

Constitutive laws

$$\varepsilon_{xx} = \frac{\sigma_{xx}}{E} - \frac{\nu}{E} \sigma_{yy}$$

$$\varepsilon_{yy} = \frac{\sigma_{yy}}{E} - \frac{\nu}{E} \sigma_{xx}$$

$$\gamma_{xy} = \frac{2(1+\nu)}{E} \tau_{xy}$$

Equilibrium equations

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + B_x = 0$$

$$\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + B_y = 0$$

Compatibility equation

$$\frac{\partial^2 \varepsilon_{xx}}{\partial y^2} + \frac{\partial^2 \varepsilon_{yy}}{\partial x^2} = \frac{\partial^2 \gamma_{xy}}{\partial x \partial y}$$



Plane stress

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- Use constitutive laws in compatibility equation – eliminate strains in compatibility using constitutive law
- Compatibility equation in terms of stress – stress compatibility equation
- Eliminate shear stress in stress compatibility equation using two equilibrium equations
- Differentiate x – direction equilibrium equation wrt 'x' – y – direction eqn. wrt 'y', add these two get shear stress
- Plug this in stress compatibility equation



Plane stress

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- The governing equation –

$$\nabla^2 (\sigma_{xx} + \sigma_{yy}) = -(1 + \nu) \left(\frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} \right)$$

This is derived making use of strain compatibility equation, constitutive laws and equilibrium equations



Summary – Plane stress and strain

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Formulation	Plane strain	Plane stress
Displacement	$\mu \nabla^2 u + (\lambda + \mu) \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + B_x = 0$ $\mu \nabla^2 v + (\lambda + \mu) \frac{\partial}{\partial y} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + B_y = 0$	$\mu \nabla^2 u + \frac{E}{2(1-\nu)} \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + B_x = 0$ $\mu \nabla^2 v + \frac{E}{2(1-\nu)} \frac{\partial}{\partial y} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + B_y = 0$
Stress	$\nabla^2 (\sigma_{xx} + \sigma_{yy}) = -\frac{1}{1-\nu} \left(\frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} \right)$	$\nabla^2 (\sigma_{xx} + \sigma_{yy}) = -(1+\nu) \left(\frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} \right)$

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Summary – Plane stress and strain

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- Basic difference – coefficients involving elastic material constants
- Solving one type of plane problem gives solution to other type also using simple transformation of elastic coefficients
- In absence of body loads – in stress formulation – governing equations of plane stress and strain are same
 - No elastic constants appear in governing equation
- Solving these equations – difficult – some text book problems



Airy's stress function

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- In stress based formulation – one equation in terms of normal stresses and two equilibrium equations for finding out the complete stress state
- Reduce the complexity of equations
- Introduce a new field variable 'Airy's stress function' (ϕ)
- Aim – reduce the governing equations from three to one
- Find the distribution of single variable ' ϕ ' – get stresses from that



Airy's stress function

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- Airy's stress function (ASF) defined for stress based formulations – both plane stress and plane strain
- Basic stress equilibrium equations –

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + B_x = 0$$

$$\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + B_y = 0$$

Body loads can be derived from potential function 'V'

Defining potential

$$\vec{B} = -\nabla V \Rightarrow B_i = V_{,i}$$

$$B_x = -\frac{\partial V}{\partial x}, \quad B_y = -\frac{\partial V}{\partial y}$$

Substitute in equilibrium equations



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Equilibrium equations become

$$\frac{\partial}{\partial x}(\sigma_{xx} - V) + \frac{\partial \tau_{xy}}{\partial y} = 0; \quad \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial}{\partial y}(\sigma_{yy} - V) = 0$$

Defining ASF as

$$\sigma_{xx} - V = \frac{\partial^2 \phi}{\partial y^2}; \quad \tau_{xy} = -\frac{\partial^2 \phi}{\partial x \partial y}; \quad \sigma_{yy} - V = \frac{\partial^2 \phi}{\partial x^2}$$

ASF, $\phi = \phi(x, y)$

ASF satisfies both equilibrium equations.

In stress based formulation – it has to satisfy one more governing equation

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- In plane strain – stress based formulation – governing equation

$$\nabla^2(\sigma_{xx} + \sigma_{yy}) = -\frac{1}{1-\nu} \left(\frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} \right)$$

$$\nabla^2 \left(\left(V + \frac{\partial^2 \phi}{\partial y^2} \right) + \left(V + \frac{\partial^2 \phi}{\partial x^2} \right) \right) = -\frac{1}{1-\nu} \left(\frac{\partial}{\partial x} \left(-\frac{\partial V}{\partial x} \right) + \frac{\partial}{\partial y} \left(-\frac{\partial V}{\partial y} \right) \right)$$

$$\nabla^2 \left(\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} \right) + 2\nabla^2 V = \frac{1}{1-\nu} \nabla^2 V$$

$$\nabla^2 \cdot \nabla^2 \phi = \left(\frac{1}{1-\nu} - 2 \right) \nabla^2 V$$

Bi-harmonic operator $\nabla^4 \phi = -\left(\frac{1-2\nu}{1-\nu} \right) \nabla^2 V$

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Airy's stress function

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- In plane stress – stress based formulation – governing equation

$$\nabla^2(\sigma_{xx} + \sigma_{yy}) = -(1 + \nu) \left(\frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} \right)$$

$$\nabla^2 \left(\left(V + \frac{\partial^2 \phi}{\partial y^2} \right) + \left(V + \frac{\partial^2 \phi}{\partial x^2} \right) \right) = -(1 + \nu) \left(\frac{\partial}{\partial x} \left(-\frac{\partial V}{\partial x} \right) + \frac{\partial}{\partial y} \left(-\frac{\partial V}{\partial y} \right) \right)$$

$$\nabla^2 \left(\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} \right) + 2\nabla^2 V = (1 + \nu) \nabla^2 V$$

$$\nabla^2 \cdot \nabla^2 \phi = (1 + \nu - 2) \nabla^2 V$$

$$\nabla^4 \phi = -(1 - \nu) \nabla^2 V$$

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Airy's stress function

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■ Governing equations in terms of ASF

$$\nabla^4 \phi = -\left(\frac{1-2\nu}{1-\nu}\right) \nabla^2 V \quad \text{Plane strain}$$

$$\nabla^4 \phi = -(1-\nu) \nabla^2 V \quad \text{Plane stress}$$

If body forces are neglected – $V = 0$ – In both cases the governing equation is same

$$\nabla^4 \phi = 0 \Rightarrow \frac{\partial^4 \phi}{\partial x^4} + 2 \frac{\partial^4 \phi}{\partial x^2 \partial y^2} + \frac{\partial^4 \phi}{\partial y^4} = 0$$

Bi-harmonic equation – solution to this equation \Rightarrow ASF

Problem of elasticity reduced to a single equation – find ϕ in solution domain and boundaries – BCs - Traction



Stress formulation

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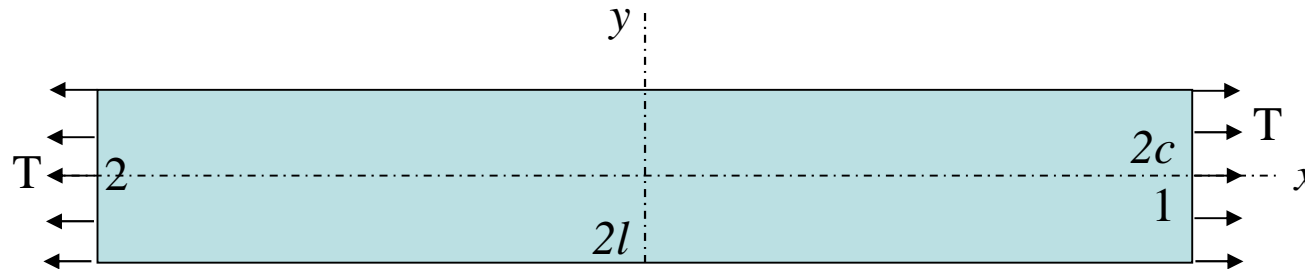
- This formulation is appropriate for use with traction boundary conditions
- Use Hooke's law for calculating strains from stresses
- Compute displacements from strain-displacement relations
- Since compatibility equations are used in formulating governing equations, the displacements obtained from integration of strain-displacement equations yield single valued, continuous field
- All equations are used for solving the problem
- Closed form analytical solutions for elasticity problems – difficult – rely on numerical methods



Uniaxial tension – ASF approach

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- A bar subjected to tension



Boundary – 1

DCs $\Rightarrow (1, 0, 0)$

$$T_x = T = l\sigma_{xx} + m\tau_{xy} = \sigma_{xx} = T$$

$$T_y = 0 = m\sigma_{yy} + l\tau_{xy} \Rightarrow \tau_{xy} = 0$$

Boundary – 2

DCs $\Rightarrow (-1, 0, 0)$

$$T_x = -T = l\sigma_{xx} + m\tau_{xy} = \sigma_{xx} = T$$

$$T_y = 0 = m\sigma_{yy} + l\tau_{xy} \Rightarrow \tau_{xy} = 0$$

Remaining boundaries – stress free.

Stress in x – direction = $T = \text{constant}$
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Uniaxial tension – ASF approach

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Stress in x – direction given by ASF

$$\sigma_{xx} = \text{constant} = \frac{\partial^2 \phi}{\partial y^2} = T \Rightarrow \phi = \frac{1}{2}Ty^2 + c_1y + c_2$$

$$\text{At } x = \pm l, \quad \sigma_{yy} \text{ and } \tau_{xy} = 0$$

$$\epsilon_x = \frac{\sigma_x}{E} - \frac{\nu}{E}\sigma_y = \frac{\partial u}{\partial x} = \frac{T}{E} \quad - \quad (1)$$

$$\epsilon_y = \frac{\sigma_y}{E} - \frac{\nu}{E}\sigma_x = \frac{\partial v}{\partial y} = -\frac{\nu T}{E} \quad - \quad (2)$$

Integrate (1) and (2)

$$u = \frac{T}{E}x + f(y)$$

$$v = -\frac{\nu T}{E}y + g(x)$$

$f(y)$ and $g(x)$ –
arbitrary functions

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Uniaxial tension – ASF approach

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Shear strain,

$$\gamma_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} = \frac{\tau_{xy}}{G} = 0$$

$$\Rightarrow f'(y) + g'(x) = 0$$

$$\Rightarrow -f'(y) = g'(x) = \omega_o$$

$$\Rightarrow \frac{\partial f}{\partial y} = -\omega_o \Rightarrow f = -\omega_o y + u_o$$

$$\frac{\partial g}{\partial x} = \omega_o \Rightarrow g = \omega_o x + v_o$$

ω_o, u_o, v_o – arbitrary constants of integration.

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Uniaxial tension – ASF approach

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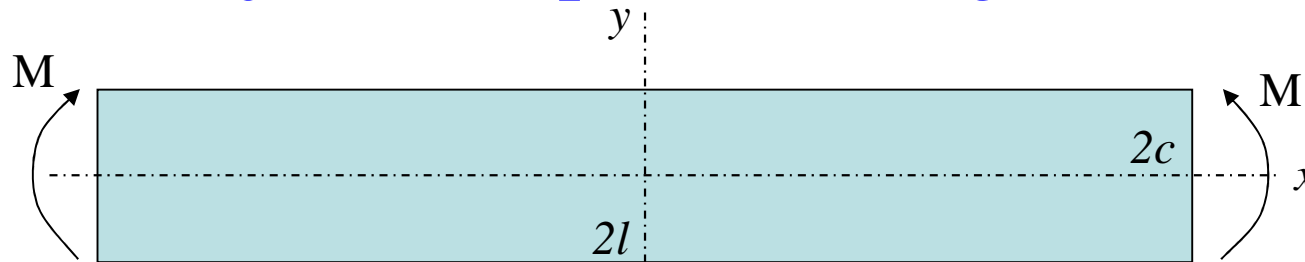
- Functions ' f ' and ' g ' represent rigid body motions
- Rigid body rotation - ω_o and translations in ' x ' and ' y ' directions $\Rightarrow u_o$ and v_o
- Terms related to rigid body motion result from strain-displacement relations
- Displacements are determined from strain fields only up to an arbitrary rigid motion
- For complete determination of displacement field – additional boundary conditions required – evaluate ω_o , u_o and v_o
- If rod has no rigid body motion $\Rightarrow \omega_o$, u_o and $v_o = 0$



Pure bending of beam – ASF approach

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- Beam subjected to pure bending –



Boundary

Conditions

$$\sigma_{yy} (x, \pm c) = 0$$

$$\tau_{xy} (x, \pm c) = 0$$

$$\tau_{xy} (\pm l, y) = 0$$

$$F_x = \int_{-c}^{+c} \sigma_{xx} (\pm l, y) dy = 0$$

Statically equivalent
boundary condition

$$-M = \int_{-c}^{+c} \sigma_{xx} (\pm l, y) y dy$$

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Pure bending of beam – ASF approach

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- Key point in ASF approach – selection of an appropriate stress function - ϕ
- It has to satisfy all boundary conditions – capture physics of the problem
- Some knowledge about structure's behavior helps in selection of ' ϕ '
 - In bending, variation of bending stress is linear along 'y' co-ordinate
 - ' ϕ ' should be a cubic polynomial



Pure bending of beam – ASF approach

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Linear variation of bending stress

$$\sigma_{xx} = Ay = \frac{\partial^2 \phi}{\partial y^2} \quad \text{Integrating the function}$$

$$\phi = \frac{A}{6} y^3 + C_1 y + C_2 \quad \text{ASF satisfies all BCs}$$

$$\int_{-c}^{+c} \sigma_{xx} (\pm l, y) y dy = -M$$

$$\int_{-c}^{+c} Ay^2 dy = -M \Rightarrow -A \left(\frac{y^3}{3} \right)_{-c}^{+c} = M \Rightarrow A = -\frac{3}{2} \frac{M}{c^3}$$

$$\phi = -\frac{3}{2} \frac{M}{c^3} \frac{y^3}{6} + C_1 y + C_2 = -\frac{M}{4c^3} y^3 + C_1 y + C_2$$

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Pure bending of beam – ASF approach

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■ Stresses and strains

$$\phi = -\frac{M}{4c^3} y^3 + C_1 y + C_2 \Rightarrow \sigma_{xx} = \frac{\partial^2 \phi}{\partial y^2} = -\frac{3M}{2c^3} y$$

$$\sigma_{yy} = 0, \quad \tau_{xy} = 0$$

$$\epsilon_{xx} = \frac{\sigma_{xx}}{E} - \nu \frac{\sigma_{yy}}{E} \Rightarrow \frac{\partial u}{\partial x} = -\frac{3}{2} \frac{M}{Ec^3} y$$

$$\epsilon_{yy} = \frac{\sigma_{yy}}{E} - \nu \frac{\sigma_{xx}}{E} = \frac{3}{2} \frac{M\nu}{Ec^3} y$$

Integrate strain –
displacement
equations

$$u = -\frac{3M}{2Ec^3} xy + f(y); \quad v = \frac{3M\nu}{4Ec^3} y^2 + g(x)$$



Pure bending of beam – ASF approach

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Shear strain

$$\gamma_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} = 0 \Rightarrow -\frac{3M}{2Ec^3}x + f'(y) + g'(x) = 0$$

$$\Rightarrow -\frac{3M}{2Ec^3}x + g'(x) = -f'(y) = \omega_o \quad \text{Integrate}$$

$$g(x) = \frac{3M}{2Ec^3} \frac{x^2}{2} + \omega_o x + v_o$$

$$f(y) = -\omega_o y + u_o$$

Functions $g(x)$ and $f(y)$ represent rigid body motions

To evaluate constants, ω_o , u_o , v_o – constraints required

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