

Electromechanical instability of large deformation in dielectric elastomers

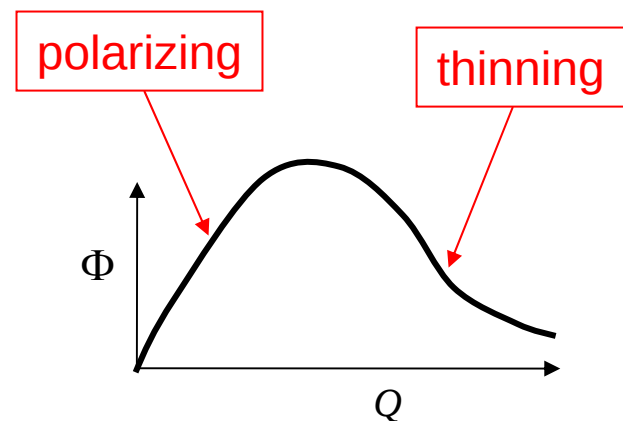
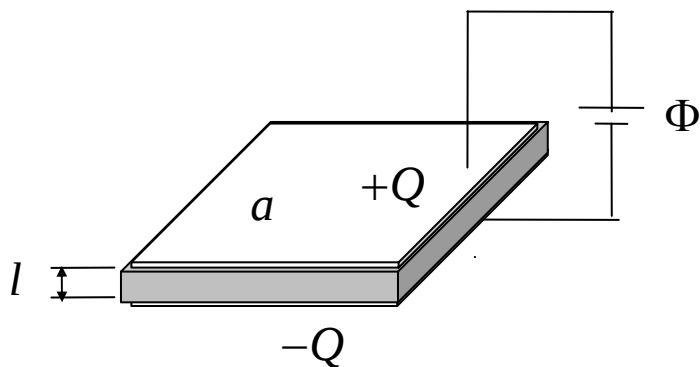
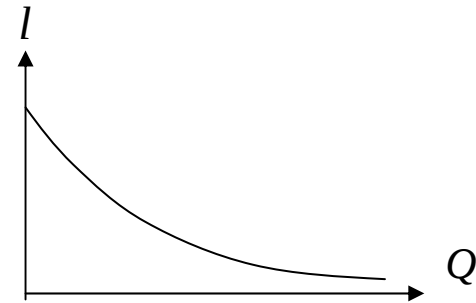
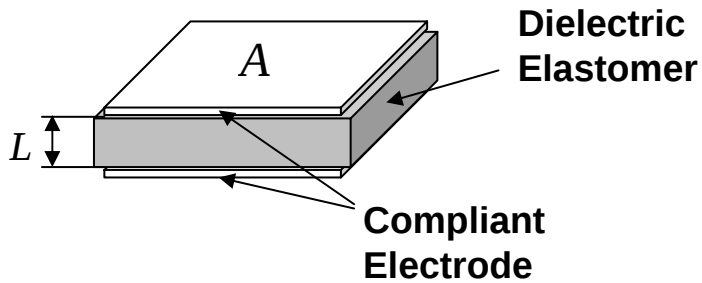
Zhigang Suo

Harvard University

Work with

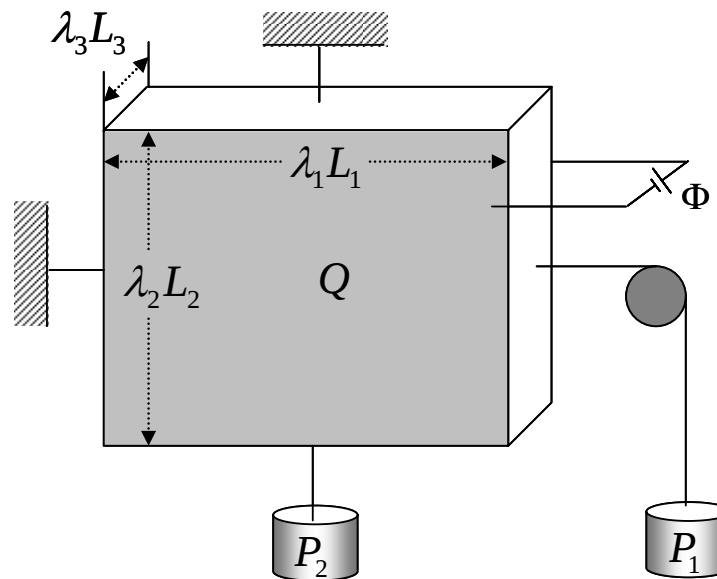
X. Zhao, W. Hong, J. Zhou, W. Greene

Pull-in instability



Theory

Zhao, Suo, APL 91, 061921 (2007)



$$s = P / A$$

$$\tilde{E} = \Phi / L$$

$$\tilde{D} = Q / A$$

Free energy
of the system

$$G(\lambda_1, \lambda_2, \tilde{D}) = L_1 L_2 L_3 W(\lambda_1, \lambda_2, \tilde{D}) - P_1 \lambda_1 L_1 - P_2 \lambda_2 L_2 - \Phi Q$$

↑ System
 ↑ Dielectric
 ↑ Weights
 ↑ Battery

Equilibrium
condition

$$0 = \frac{\delta G}{L_1 L_2 L_3} = \left(\frac{\partial W}{\partial \lambda_1} - s_1 \right) \delta \lambda_1 + \left(\frac{\partial W}{\partial \lambda_2} - s_2 \right) \delta \lambda_2 + \left(\frac{\partial W}{\partial \tilde{D}} - \tilde{E} \right) \delta \tilde{D}$$

Need free-energy function $W(\lambda_1, \lambda_2, \tilde{D})$

Dielectric constant is insensitive to deformation

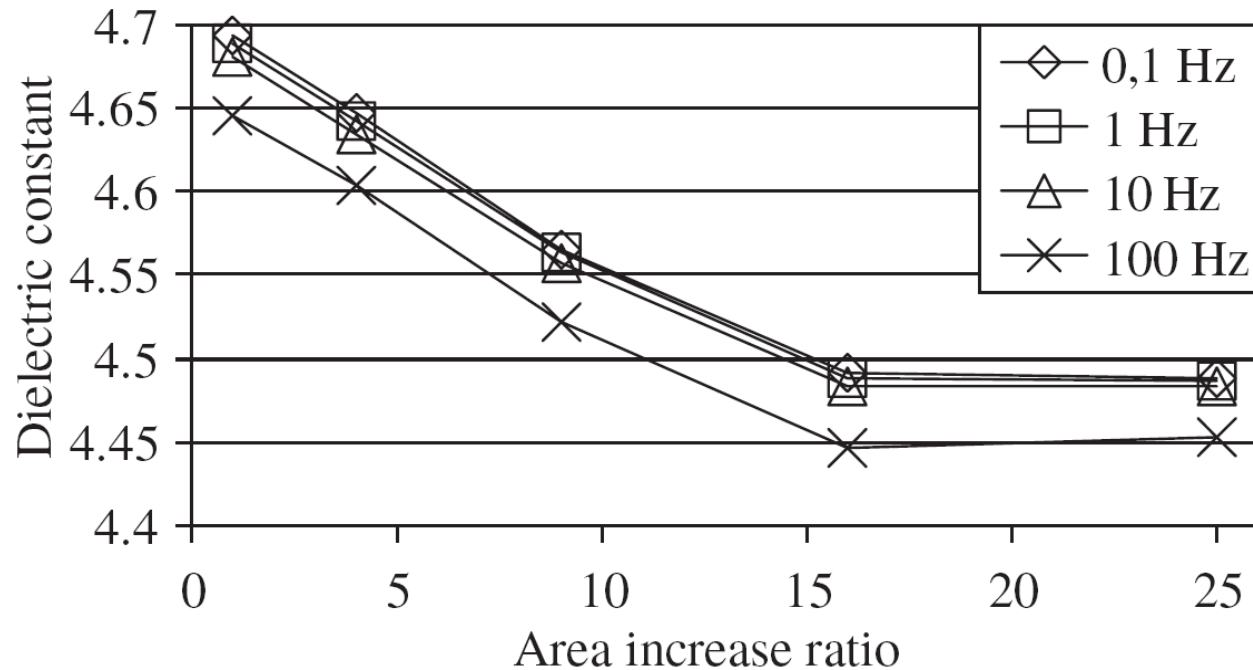


Figure 5. The relative dielectric constant of VHBTM 4910 drops, when it is stretched.

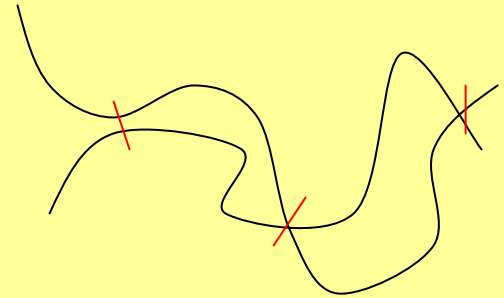
Material law: Ideal dielectric elastomers

Zhao, Hong, Suo, Physical Review B 76, 134113 (2007)

Dielectric behavior is liquid-like, unaffected by deformation.

$$W(\lambda_1, \lambda_2, \tilde{D}) = W_s(\lambda_1, \lambda_2) + \frac{D^2}{2\epsilon}$$

↑ ↑
Stretch Polarization



Ideal electromechanical coupling is purely a geometric effect:

$$D = \frac{Q}{a}$$

$$\tilde{D} = \frac{Q}{A}$$

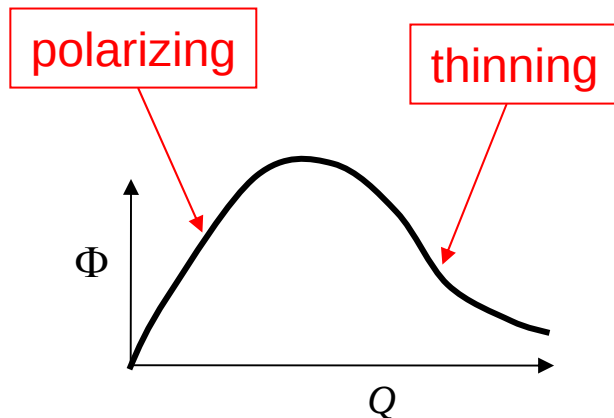
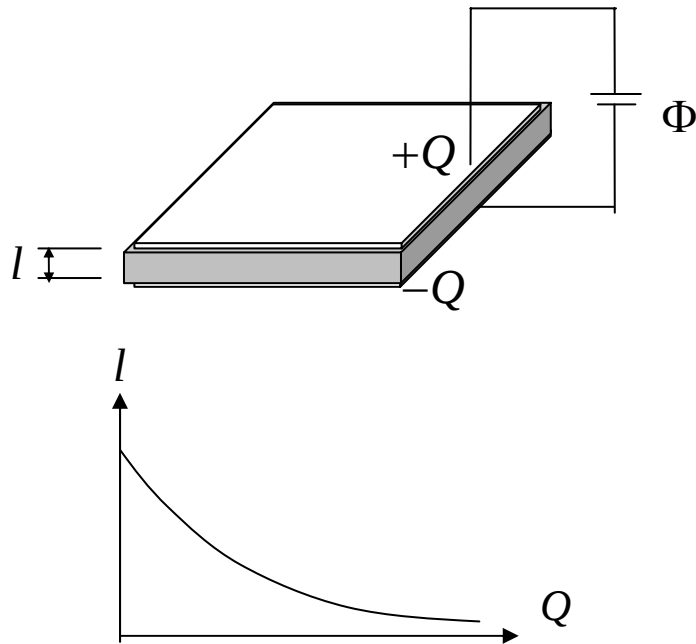
$$D = \frac{\tilde{D}}{\lambda_1 \lambda_2}$$

Choose a free energy of stretching. For example, neo-Hookean law:

$$W_s(\lambda_1, \lambda_2) = \frac{\mu}{2} (\lambda_1^2 + \lambda_2^2 + \lambda_3^2 - 3)$$

incompressibility $\lambda_3 = 1/(\lambda_1 \lambda_2)$

Theory of pull-in instability



$$W(\lambda, \tilde{D}) = \frac{\mu}{2}(\lambda^2 + 2\lambda - 3) + \frac{\lambda^2 \tilde{D}^2}{\varepsilon}$$

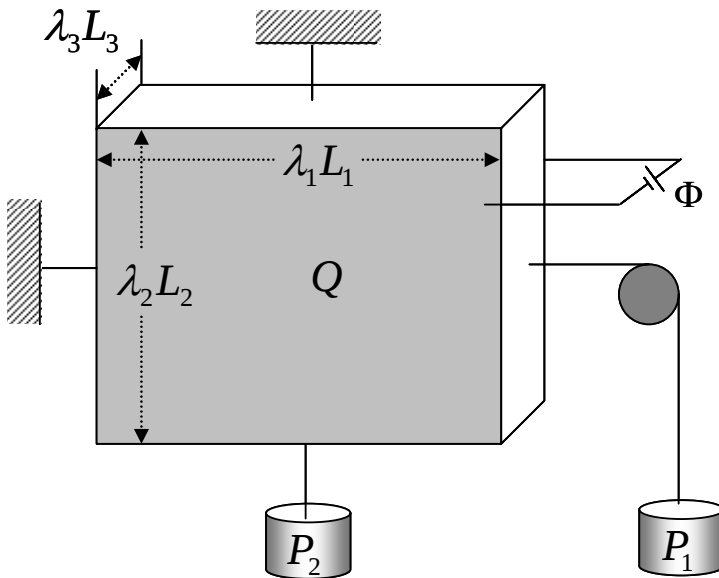
$$s = \frac{\partial W(\lambda, \tilde{D})}{\partial \lambda} = 0 \quad \tilde{E} = \frac{\partial W(\lambda, \tilde{D})}{\partial \tilde{D}}$$

$$\lambda = \left(1 + \frac{\tilde{D}^2}{\varepsilon \mu} \right)^{-1/3}$$

$$\frac{\tilde{E}}{\sqrt{\mu / \varepsilon}} = \frac{\tilde{D}}{\sqrt{\varepsilon \mu}} \left(1 + \frac{\tilde{D}^2}{\varepsilon \mu} \right)^{-2/3}$$

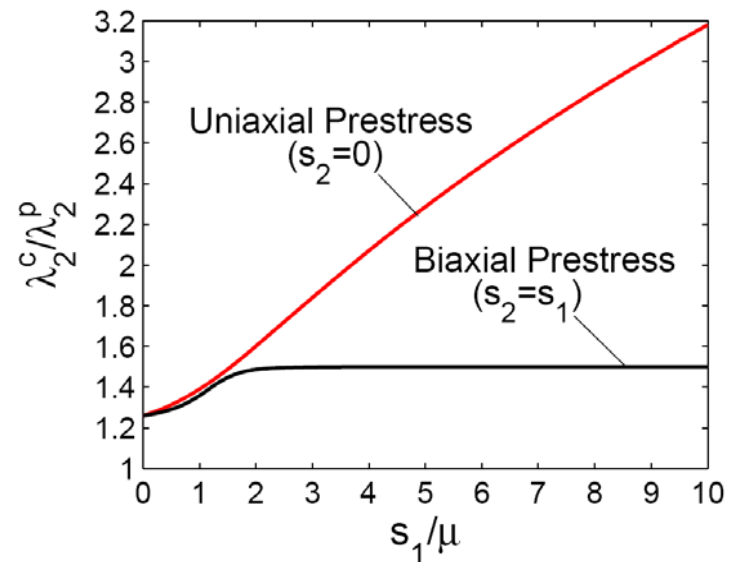
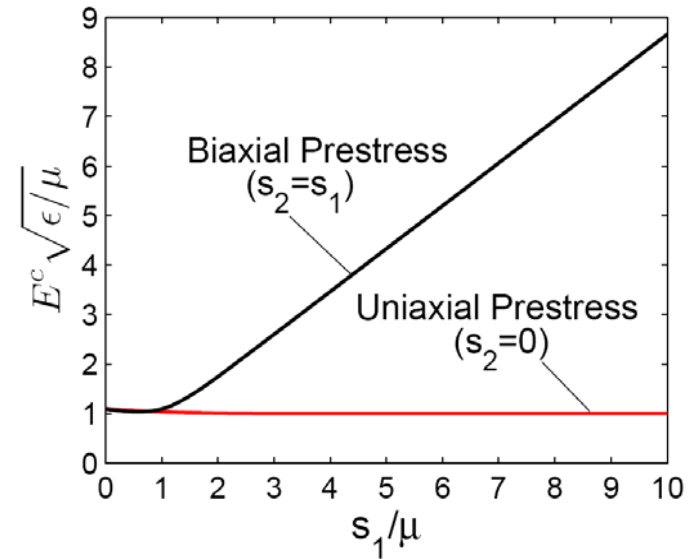
$$\lambda_c \approx 0.63 \quad \tilde{E}_c \sim \sqrt{\frac{\mu}{\varepsilon}} \sim \sqrt{\frac{10^6 \text{ N/m}}{10^{-10} \text{ F/m}}} = 10^8 \text{ V/m}$$

Pre-stretch increases actuation strain

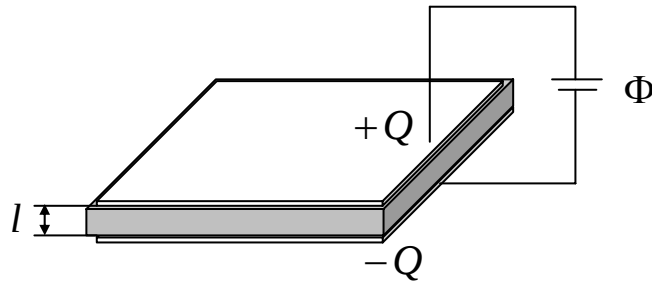


Experiment: Pelrine, Kornbluh, Pei, Joseph
Science 287, 836 (2000).

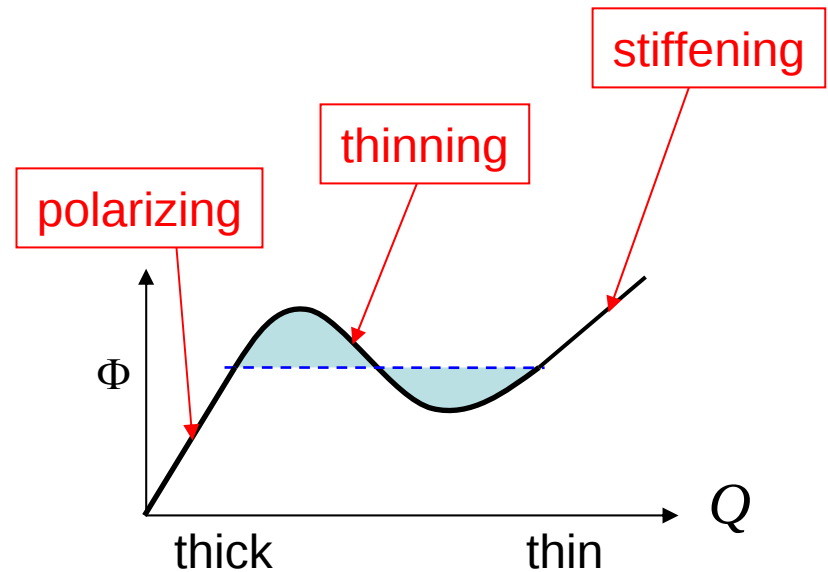
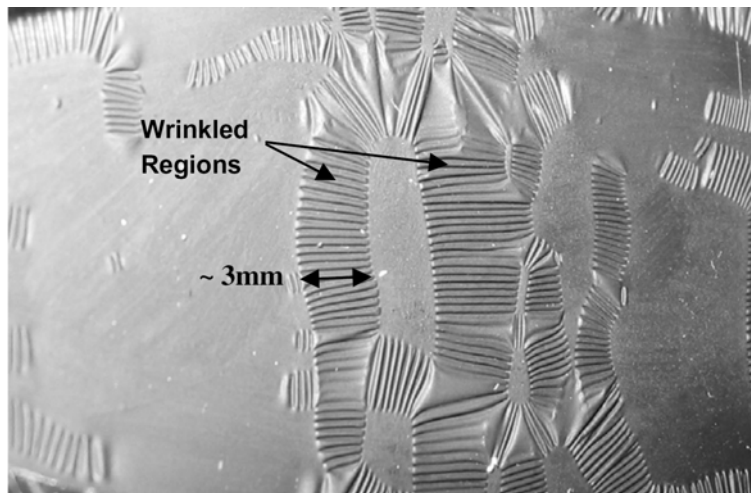
Theory: Zhao, Suo
APL 91, 061921 (2007)



Coexistent states



Top view



Cross section

Coexistent states: flat and wrinkled

Experiment: Plante, Dubowsky,
Int. J. Solids and Structures **43**, 7727 (2006)

Theory: Zhao, Hong, Suo
Physical Review B **76**, 134113 (2007)

Stiffening due to extension limit

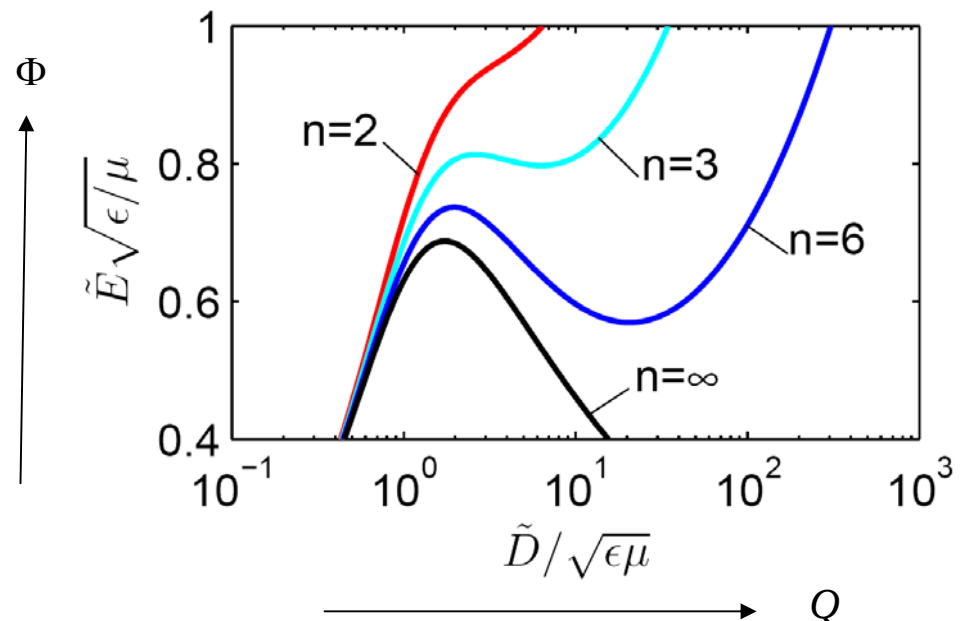
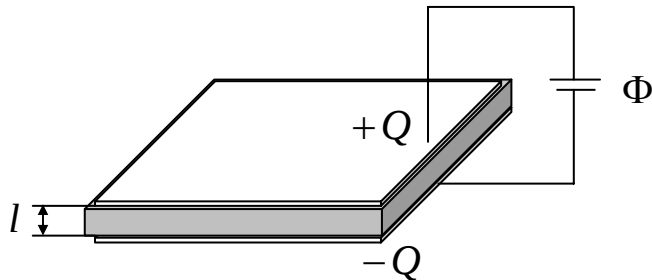
Stiffening as each polymer chain approaches its fully stretched length
(e.g., Arruda-Boyce model)

$$W_s = \mu \left[\frac{1}{2} (I - 3) + \frac{1}{20n} (I^2 - 9) + \dots \right]$$

$$I = \lambda_1^2 + \lambda_2^2 + \lambda_3^2$$

μ : small-strain shear modulus

n : number of monomers per chain



FEM: inhomogeneous deformation



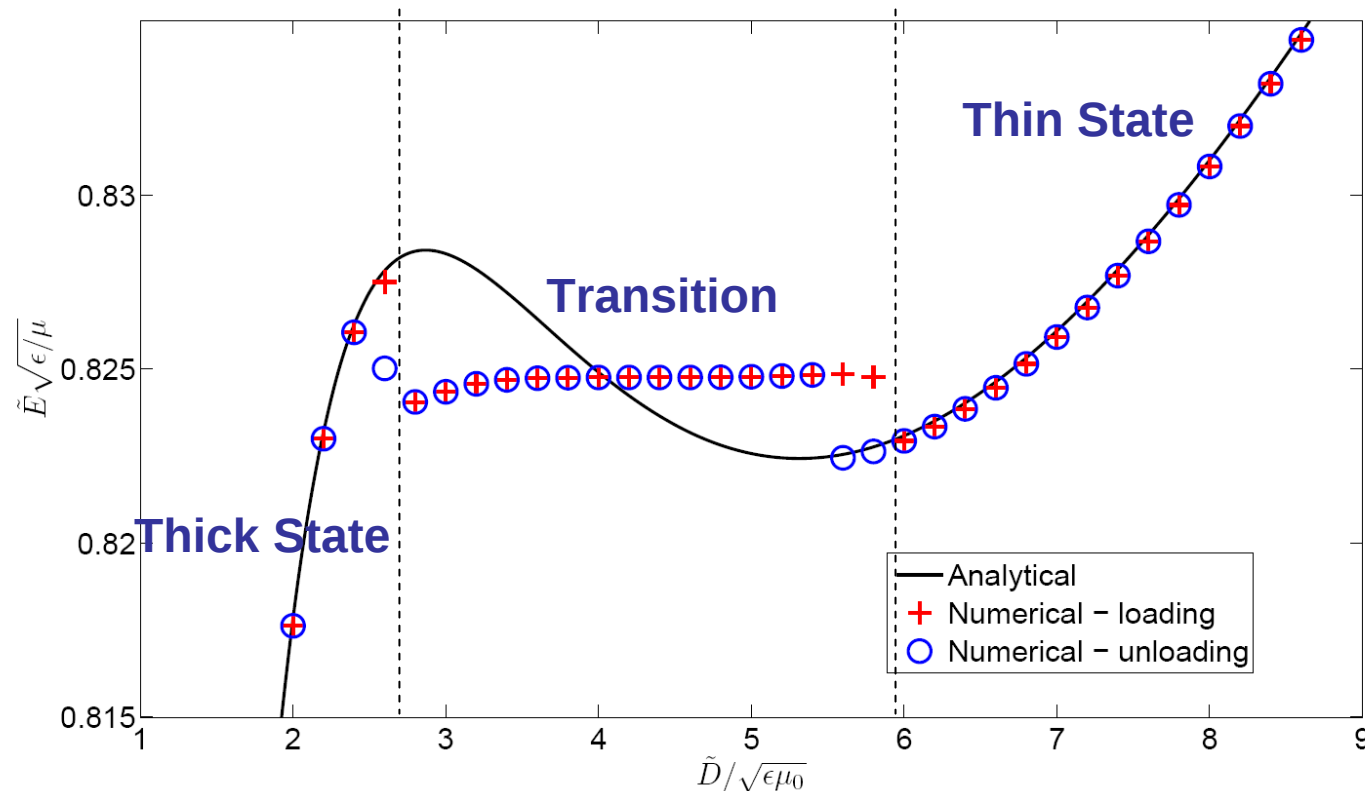
Thick State



Transition



Thin State



Summary

- **Ideal dielectric elastomers:** dielectric behavior is liquid like, unaffected by deformation.
- **Pull-in instability:** coupling large deformation and electric charge.
- **Pre-stretch** increases actuation strain.
- **Co-existent states:** stiffening due to extension limit.

Game plan

- Inhomogeneous deformation.
- Quantitative comparison with experiments.
- Effect of viscosity on instability.
- Nonideal electromechanical coupling.
- Adding other effects (solvents, ions, enzymes...): Soft Active Materials (SAMs)

<http://imechanica.org/node/2280>